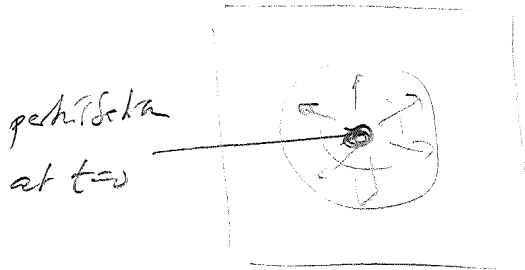


How does information propagate in spin systems?



"speed of sound"?

Relativistic systems \rightarrow speed of light.

Non-relativistic \rightarrow a priori no such limit.

WIM sees For spin systems on lattice, such limit exists!

Why interesting?

- Fundamental relevance: finite propagation speed are w/out relativity
- Needed to define dynamics in the thermodyn. limit
- allows to transfer a range of "relativistic" results to non-rel. setting

Definitions.

Lattice:

- Consider general lattice: $\Lambda = \text{set of lattice points}$
- Consider finite (but very big) lattice
- notion of distance (metric) $d(i,j): \Lambda \times \Lambda \rightarrow \mathbb{R}_{\geq 0}$.
(and Δ req.: $d(i,j) \leq d(i,k) + d(k,j)$)

E.g.: Square Lattice: $\Lambda = \{i(i_1, i_2) : 0 \leq i_1, i_2 \leq N-1\}$

$d(i, j)$: e.g. Euclidean distance, or "Manhattan metric", possibly w/ periodic bnds.

Define: For $X, Y \subset \Lambda$: $d(X, Y) = \min_{\substack{i \in X \\ j \in Y}} d(i, j)$

$\text{diam}(X) = \max_{i, j \in X} d(i, j)$

$|X| = \#$ of elements in X

Hilbert space: To each lattice site, associate $\mathcal{H}_i = \mathbb{C}^d$.

$\mathcal{H}_X = \bigotimes_{i \in X} \mathcal{H}_i$; $\mathcal{H}_\Lambda =$ total H.S. of system

Operators: $O: \mathcal{H}_\Lambda \rightarrow \mathcal{H}_\Lambda$.

If $O = A_X \otimes \mathbb{1}_{\Lambda \setminus X}$, we say O is supported on X , $X = \text{supp } O$.

Operator norm: $\|O\| = \max_{|\psi\rangle} \frac{|O|\psi\rangle|}{\|\psi\rangle}$.

$\|A \otimes B\| = \|A\| \|B\|$; $\|\mathbb{1}\| = 1$; $\|A \cdot B\| \leq \|A\| \|B\|$

$$H = \sum_{z \in \Lambda} h_z, \text{ with } \text{supp } h_z = z.$$

$$H_X = \sum_{z \cap X \neq \emptyset} h_z, \quad H \equiv H_\Lambda.$$

We will assume that H is local or rapidly decaying, i.e.,

$\|h_z\|$ is zero/small if $\text{diam}(z)$ gets large.

Theorem (Lieb-Robinson bound)

Let H be such that there exist constants $\mu, S \geq 0$ s.t. for all $i \in \Lambda$

$$\sum_{X \ni i} \|h_X\| |X| e^{\mu \cdot \text{diam}(X)} \leq S, \quad (1)$$

Then, for any A_X, B_Y w. $\text{supp } A_X = X, \text{supp } B_Y = Y, d(X, Y) > 0$, it holds that

$$\|[A_X(t), B_Y]\| \leq 2 \|A_X\| \|B_Y\| |X| e^{-\mu d(X, Y)} (e^{2St} - 1). \quad (2)$$

Here, $O(t) = e^{iHt} O e^{-iHt}$ is the time-evolved op. in the

Heisenberg picture.

Comments:

① What does (1) have to do with the locality of H ?

- o H local, e.g. 2-body Ham. w/ $\|h_x\| = \eta$ for all 2-body terms, (14)
and 2D square lattice:

$$\sum_{\substack{x \ni i \\ 4x}} \underbrace{\|h_x\|}_{=\eta} \underbrace{|x|}_{=2} e^{\mu \cdot \frac{\text{diam}(x)}{=1}} = \underline{\underline{8\eta e^\mu =: s}}$$

→ Different pairs μ & s possible.

- o $\|h_x\| \sim e^{-\gamma \text{diam}(x)} \Rightarrow$ choose $\mu \ll \gamma$ & s accordingly

② What does (2) have to do w/ speed limit?

$$e^{-\mu d(x,y)} \left(e^{2s|t|} - 1 \right) \leq e^{-\mu \left(\underbrace{d(x,y)}_{\text{distance}} - \underbrace{\frac{2s}{\mu}|t|}_{\text{velocity}} \right)}$$

prefactor.

By choosing μ we can optimize velocity. Above:

$$\frac{16\eta e^\mu}{\mu} \text{ minimal for } \mu = 1 \Rightarrow \underline{\underline{v_{LE} = 16 \cdot e \cdot \eta}}$$

③ What does the commutator "mean"?

$$[A(t), B] = 0 \iff A(t) \text{ is only supported outside supp } B,$$

Recall:

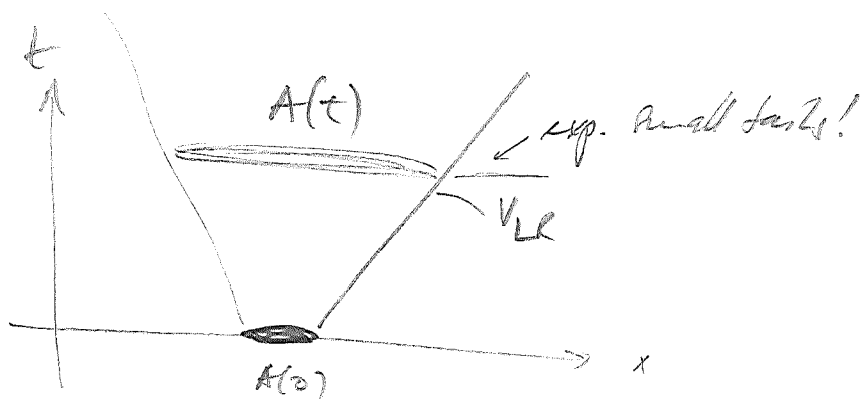
(15)

Let $K_e(x) = \{i \mid d(i, x) \leq e\}$, and B w/ τ_{pp} on $A \setminus K_e(x)$.

$[A(t), B] = 0 \iff A(t)$ supported on $K_e(x) \implies$ effect of A cannot be seen outside that e .

This statement "epilocalizes". i.e.: $d(x, y) - \frac{2s}{\mu} |t| < e$

$\implies A(t)$ localized in region $\frac{2s}{\mu} |t|$.



Note: Very important that RHS of (2) indep. of |t|!

Proof: 3 steps

Step 1: Express $\|[A(t), B]\|$ in terms of $\|[h_z(s), B]\|$.

Step 2: Eliminate $A(t)$ by defining "worst case" commutator

$$C_B(x, t) = \sup_{A \in \mathcal{A}_x} \frac{\|[A(t), B]\|}{\|A\|}$$

\uparrow algebra of all ops. w/ supp. on X .

Recursive expression for $G_B(X, t) \Rightarrow$ series for $G_B(X, t)$ in terms of $\|h_2\|$. (6)

Step 3: Bound necessary decay of $\|h_2\|$ - not used in Steps 1, 2!

Step 1: ($t > 0$):

$$\| [A(t), B] \| - \| [A(0), B] \| = \sum_{u=0}^{N-1} \epsilon \frac{\| [A(t_{u+1}), B] \| - \| [A(t_u), B] \|}{\epsilon} \quad (*)$$

with $\epsilon = \frac{t}{N}$, $t_u = \frac{t}{N} u = \epsilon u$.

Now use that (i) $\|u^\dagger 0 u\| = \|0\|$, and (ii) $A(\epsilon) = A + i\epsilon [H_x, A] + O(\epsilon^2) =$

$$\| [A(t_{u+1}), B] \| - \| [A(t_u), B] \| \stackrel{(i)}{=} \| [A(\epsilon), B(-t_u)] \| - \| [A, B(-t_u)] \|$$

$\rightarrow (u/u = \exp(iH_x t_u))$

$$= \| [A + i\epsilon [H_x, A] + O(\epsilon^2), B(-t_u)] \| - \| [A, B(-t_u)] \|$$

Δ -neg. \uparrow thus commute!

$$= \| [A + i\epsilon [H_x, A], B(-t_u)] \| - \| [A, B(-t_u)] \| + O(\epsilon^2)$$

We can now use $A + i\epsilon [H_x, A] = e^{i\epsilon H_x} A e^{-i\epsilon H_x} + O(\epsilon^2)$,

$$\| [A + i\epsilon [H_x, A], B(-t_u)] \| \leq \| [e^{i\epsilon H_x} A e^{-i\epsilon H_x}, B(-t_u)] \| + O(\epsilon^2)$$

$$= \| [A, e^{-i\epsilon H_x} B(-t_u) e^{i\epsilon H_x}] \| + O(\epsilon^2)$$

$$\leq \| [A, B(-t_u) - i\varepsilon [H_x, B(-t_u)]] \| + O(\varepsilon^2)$$

(17)

$$\stackrel{\Delta\text{-ineq.}}{\leq} \| [A, B(-t_u)] \| + \varepsilon \| [A, [H_x, B(-t_u)]] \| + O(\varepsilon^2).$$

in RHS of (**),

$$\| [A(t_{uH}), B] \| - \| [A(t_u), B] \| \leq \varepsilon \| [A, [H_x, B(-t_u)]] \| + O(\varepsilon^2).$$

$$\leq 2\varepsilon \| A \| \cdot \| [H_x, B(-t_u)] \| + O(\varepsilon^2)$$

in ~~(*)~~:

$$\| [A(\varepsilon), B] \| - \| [A(0), B] \| \leq 2 \| A \| \sum_{u=0}^{N-1} \varepsilon \cdot \| [H_x(t_u), B] \| + O(\varepsilon)$$

$$H_x = \sum_{z \in \mathbb{Z}^n \setminus \phi} h_x$$

$$\leq 2 \| A \| \sum_{z \in \mathbb{Z}^n \setminus \phi} \sum_{u=0}^{N-1} \varepsilon \| [h_z(t_u), B] \| + O(\varepsilon).$$

Limit $\varepsilon \rightarrow 0$ (as fuck by the way):

$$\| [A(t), B] \| - \| [A(0), B] \| \leq 2 \| A \| \sum_{z \in \mathbb{Z}^n \setminus \phi} \int_0^t ds \| [h_z(s), B] \|.$$

(***)

→ We have re-expressed $\| [A(t), B] \|$ w/out using A (except for $\| A \|$).

Step 2: Define 'worst case concentration' (normalised!)

(18)

$$C_B(x,t) := \sup_{A \in \mathcal{A}_x} \frac{\| [A(t), B] \|}{\|A\|}$$

(4*) $\Rightarrow C_B(x,t) \leq C_B(x,0) + 2 \sum_{z \cap X \neq \emptyset} \|h_z\| \int_0^t ds C_B(z,s).$

(4*)

Will now solve (4*) iteratively.

Initial conditions: $C_B(z,0) = 0$ if $z \cap Y = \emptyset$

$$C_B(z,0) \leq 2\|B\| \text{ if } z \cap Y \neq \emptyset.$$

In particular: $X \cap Y = \emptyset \Rightarrow C(x,0) = 0.$

Iterative solution of (4*):

$$C_B(x,t) \leq \overbrace{C_B(x,0)}^{=0} + 2 \sum_{z \cap X \neq \emptyset} \|h_z\| \int_0^t ds_1 C_B(z,s_1)$$

$$\leq 2 \cdot \sum_{z_1 \cap X \neq \emptyset} \|h_{z_1}\| \cdot \int_0^t ds_1 \cdot C_B(z_1,0)$$

$$+ 2^2 \cdot \sum_{z_1 \cap X \neq \emptyset} \|h_{z_1}\| \sum_{z_2 \cap z_1 \neq \emptyset} \|h_{z_2}\| \int_0^t ds_1 \int_0^{s_1} ds_2 C_B(z_2,s_2)$$

$$\leq \dots \text{ etc. } \leq$$

NR: Use that $C(z, 0) = \begin{cases} 2\|B\|, & z \cap Y \neq \emptyset \\ 0, & \text{else} \end{cases}$

$$\rightarrow \sum_{z_1} \int_0^t C_{\beta}(z_1, 0) = \sum_{z_1 \cap Y \neq \emptyset} t \cdot 2\|B\|$$

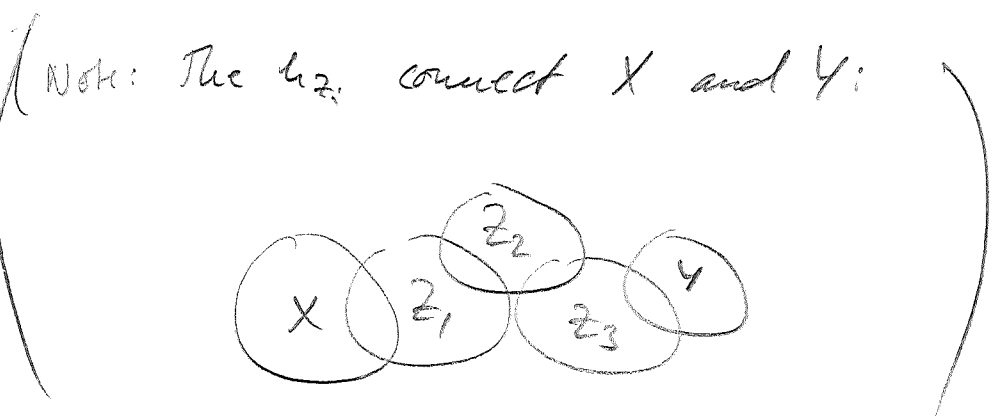
$$\sum_{z_2} \int_0^t \int_0^{t_1} C_{\beta}(z_2, 0) = \sum_{z_2 \cap Y \neq \emptyset} \frac{t^2}{2!} \cdot 2\|B\|, \text{ etc.}$$

$$\dots \leq 2\|B\| \frac{2t}{1!} \sum_{\substack{z_1 \cap X \neq \emptyset \\ z_1 \cap Y \neq \emptyset}} \|h_{z_1}\| +$$

$$+ 2\|B\| \frac{(2t)^2}{2!} \sum_{z_1 \cap X \neq \emptyset} \|h_{z_1}\| \sum_{\substack{z_2 \cap z_1 \neq \emptyset \\ z_2 \cap Y \neq \emptyset}} \|h_{z_2}\|$$

$$+ 2\|B\| \frac{(2t)^3}{3!} \sum_{z_1 \cap X \neq \emptyset} \|h_{z_1}\| \sum_{z_2 \cap z_1 \neq \emptyset} \|h_{z_2}\| \sum_{\substack{z_3 \cap z_2 \neq \emptyset \\ z_3 \cap Y \neq \emptyset}} \|h_{z_3}\|$$

+ ...



Step 3: Bound the sum using the decay of the $\|h_z\|$. (20)

(NB: For strictly local $\|h_z\|$, this can be done using results from percolation theory, as $\sum_{\ell} \sum \dots = \neq \text{paths} \cdot \|h\|(\ell)$)

Term by term:

1st term:

$$\sum_{\substack{z_1 \cap X \neq \emptyset \\ z_1 \cap Y \neq \emptyset}} \|h_{z_1}\| \leq \sum_{i \in X} \sum_{\substack{z_1 \ni i \\ z_1 \cap Y \neq \emptyset}} \|h_{z_1}\| e^{\underbrace{\mu d(i, Y)}_{\leq \text{diam}(z_1)} - \mu d(i, Y)} = 1$$

$$\leq \sum_{i \in X} \left(\underbrace{\sum_{\substack{z_1 \ni i \\ z_1 \cap Y \neq \emptyset}} \|h_{z_1}\| e^{\mu \text{diam}(z_1)} \cdot \frac{1}{|z_1|}}_{\text{decay} \leq S} \right) e^{-\mu d(i, Y)}$$

$$\leq S \sum_{i \in X} e^{-\mu d(i, Y)}$$

$$\Rightarrow \text{1st term} = 2 \|B\| \frac{(2ts)^t}{t!} \sum_{i \in X} e^{-\mu d(i, Y)}$$

2nd term: Use $d(i, Y) \leq d(i, j) + d(j, Y)$ (Δ -Ineq.)

$$\Rightarrow e^{-\mu d(i, Y)} e^{\mu d(i, j)} e^{\mu d(j, Y)} \geq 1$$

Then:

$$\sum_{z_1 \cap X \neq \emptyset} \|h_{z_1}\| \sum_{\substack{z_2 \cap z_1 \neq \emptyset \\ z_2 \cap Y \neq \emptyset}} \|h_{z_2}\| \leq$$

$$\leq \sum_{i \in X} \sum_{z_1 \ni i} \|h_{z_1}\| \sum_{\substack{j \in z_1 \\ z_2 \ni j \\ z_2 \cap Y \neq \emptyset}} \|h_{z_2}\| \leq$$

$$\leq \sum_{i \in X} e^{-\mu d(i, Y)} \sum_{z_1 \ni i} \|h_{z_1}\| \sum_{j \in z_1} e^{\mu d(i, j)} \underbrace{\sum_{\substack{z_2 \ni j \\ z_2 \cap Y \neq \emptyset}} \|h_{z_2}\| e^{\mu d(j, Y)}}_{\leq e^{\mu \text{diam}(z_2)}} \leq S$$

$$\leq S \sum_{i \in X} e^{-\mu d(i, Y)} \underbrace{\sum_{z_1 \ni i} \|h_{z_1}\| \sum_{j \in z_1} e^{\mu d(i, j)}}_{\leq e^{\mu \text{diam}(z_1)}} \leq |z_1| e^{\mu \text{diam}(z_1)} \leq S$$

$$\leq S^2 \sum_{i \in X} e^{-\mu d(i, Y)}$$

$$\Rightarrow \text{2nd term} \geq 2 \|B\| \frac{(2ts)^2}{2!} \sum_{i \in X} e^{-\mu d(i, Y)}$$

etc. ...

Together:

$$C_B(x,t) \leq 2 \|B\| \sum_{k \geq 1} \frac{(2ts)^k}{k!} \sum_{i \in X} e^{-\mu d(i,y)}$$

$$= 2 \|B\| (e^{2ts} - 1) \sum_{i \in X} e^{-\mu d(i,y)}$$

$$\boxed{d(i,y) \geq \text{dist}(x,y)}$$

$$\leq 2 \|B\| (e^{2ts} - 1) e^{-\mu \text{dist}(x,y)} \cdot |X|$$

□